

# Ph 406: Elementary Particle Physics

## Problem Set 8

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1. The spin-1 vector mesons can be taken to have quark content:  $\rho^0 = (u\bar{u} - d\bar{d})/\sqrt{2}$ ,  $\omega^0 = (u\bar{u} + d\bar{d})/\sqrt{2}$ ,  $\phi = s\bar{s}$ ,  $J/\psi = c\bar{c}$ ,  $\Upsilon = b\bar{b}$  ( $V_{\text{top}} = t\bar{t}$  will not exist).

The decays  $V \rightarrow e^+e^-$  proceed via a single intermediate photon, where  $V$  is a vector meson. In the quark model, this corresponds to the reaction  $q\bar{q} \rightarrow \gamma \rightarrow e^+e^-$ , whose cross section was discussed on p. 108, Lecture 7 of the Notes. Deduce the decay rate  $\Gamma$  for this by recalling (p. 13, Lecture 1 of the Notes) that

$$\text{Rate} = \Gamma_{a+b \rightarrow c+d} = N v_{\text{rel}} \sigma_{a+b \rightarrow c+d}, \quad (1)$$

where  $N$  is the number of candidate scatters per second per unit volume, and  $v_{\text{rel}}$  is the relative velocity of the initial-state particles  $a$  and  $b$ . In case of a two-particle bound state,  $N = |\psi(0)|^2$  is the probability that both particles are at the origin.

Predict the decay rates to  $e^+e^-$  for the five vector mesons in the model that the strong interaction between (colored) quarks at short distances can be described by the Coulomb-like potential  $V(r) \approx -4\alpha_S/3r$ .

Compare with data summarized at [http://pdg.lbl.gov/2013/tables/contents\\_tables\\_mesons.html](http://pdg.lbl.gov/2013/tables/contents_tables_mesons.html).

What do we learn from this about possible energy dependence of  $\alpha_S$ ?

2. In the vector-meson decays  $V \rightarrow \pi^0\gamma$ ,  $\eta\gamma$ , the meson spin changes from 1 to 0. Hence, this must be an M1 (magnetic dipole) transition. In the quark model the M1 electromagnetic transition flips a single quark spin, but does not change quark flavor, with matrix element proportional to the relevant quark magnetic moment(s). Suppose the quarks have Dirac moments  $Q_q/2m_q$  where  $m_u \approx m_d \approx \frac{2}{3}m_s$ . Predict the relative decay rates (not just matrix elements) to  $\pi^0\gamma$  and  $\eta\gamma$  for the  $\rho^0$ ,  $\omega^0$ ,  $\phi$  and  $J/\psi$  vector mesons.

Recall that in the quark model the spin-0-octet neutral mesons have quark wavefunction  $\pi^0 = (u\bar{u} - d\bar{d})/\sqrt{2}$  and  $\eta(548) = (2s\bar{s} - u\bar{u} - d\bar{d})/\sqrt{6}$ . Compare with data summarized at [http://pdg.lbl.gov/2013/tables/contents\\_tables\\_mesons.html](http://pdg.lbl.gov/2013/tables/contents_tables_mesons.html).

3. The  $\psi'(3685)$  vector meson can decay to  $\chi(3415) + \gamma$ . The  $\chi$  particle is believed to be a  $^3P_0$   $c\bar{c}$  state. If so, predict the angular distribution of the  $\gamma$  relative to the direction of the electron supposing the  $\psi'$  is produced in a colliding-beam experiment  $e^+e^- \rightarrow \psi' \rightarrow \chi\gamma$ . Recall that at high energies the one-photon annihilation of  $e^+e^-$  proceeds entirely via transversely polarized photons ( $S_z = \pm 1$ ).

## Crossing Symmetry

We have previously noted that the inverse processes  $a + b \leftrightarrow c + d$  have common matrix elements, and that these process may proceed via single-particle exchange in any of the  $s$ -,  $t$ - or  $u$ -channels with related matrix elements. Such relations among matrix elements for related processes are sometimes called **crossing symmetry**.

In the next 3 problems you will use crossing symmetry to convert the matrix element for muon decay,  $\mu \rightarrow e\bar{\nu}_e\nu_\mu$  to results for 3 related processes.

The square of the matrix element for the 4-particle vertex  $\mu\nu_\mu e\nu_e$  of unpolarized particles in the Fermi theory of the weak interaction is, in terms of the particle 4-vectors  $p$ ,

$$|M|^2 = 32G_F^2(p_\mu \cdot p_{\nu_e})(p_e \cdot p_{\nu_\mu}), \quad (2)$$

where  $G_F$  is Fermi's constant, and the average over initial spins and sum over final spins is the same for all variants of the vertex.

4. Deduce the cross section for the neutrino-scattering reaction  $\nu_\mu + e^- \rightarrow \mu^- + \nu_e$ .

Recall that the differential cross section for 2-particle scattering  $a + b \rightarrow c + d$  can be written in the center of mass frame as (p. 80, Lecture 5 of the Notes)

$$\frac{d\sigma}{d\Omega^*} = \frac{|M|^2 P_f}{64\pi^2 s P_i}, \quad (3)$$

where  $P_f$  is the momentum of the final-state particles  $c$  and  $d$ ,  $P_i$  is the momentum of the initial-state particles  $a$  and  $b$ , and  $s = (p_a + p_b)^2 = (p_c + p_d)^2$  is the square of the total energy in the center of mass frame. Express the cross section in terms of  $s$ , and then evaluate this in the lab frame where the electron is at rest and the muon neutrino has energy  $E$ .

5. The process  $\mu^+e^- \rightarrow \mu^-e^+$  was considered by Pontecorvo in 1957 as a possible example of quantum oscillations of a two-particle system,<sup>1</sup> as this reaction could proceed via a two-neutrino intermediate state.

While the reaction  $\mu^+e^- \rightarrow \bar{\nu}_\mu\nu_e$  is unlikely ever to be observed, it is now understood that the related reaction  $e^+e^- \rightarrow \bar{\nu}_e\nu_e$  is the main source of neutrino production in supernovae, and a key process in their history.

Use a suitable variant of the matrix element (2) to deduce the cross section for  $e^+e^- \rightarrow \bar{\nu}_e\nu_e$ . *Work in the center-of-mass frame, and express the result in terms of the invariant  $s$ .*

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<sup>1</sup>B. Pontecorvo, *Mesonium and Antimesonium*, Sov. Phys. JETP **6**, 429 (1957),

[http://kirkmcd.princeton.edu/examples/EP/pontecorvo\\_spjjetp\\_6\\_429\\_57.pdf](http://kirkmcd.princeton.edu/examples/EP/pontecorvo_spjjetp_6_429_57.pdf).

This landmark paper introduced the term mesonium, raised the possibility that  $\nu_e$  and  $\nu_\mu$  are different particles, made the first speculations about neutrino oscillations, and led to the notion of conservation of lepton number, as developed by G. Feinberg and S. Weinberg, *Law of Conservation of Muons*, Phys. Rev. Lett. **6**, 381 (1961), [http://kirkmcd.princeton.edu/examples/EP/feinberg\\_prl\\_6\\_381\\_61.pdf](http://kirkmcd.princeton.edu/examples/EP/feinberg_prl_6_381_61.pdf).

The divergence of this cross section at low energy is avoided by Nature in that the electron and positron would not scatter but rather would bind into a positronium atom.

6. Some positronium atoms (which of the ortho- and para- states?) can decay to two neutrinos. Deduce the decay rate  $\Gamma$  for this by recalling (p. 13, Lecture 1 of the Notes) that

$$\text{Rate} = \Gamma_{a+b \rightarrow c+d} = N v_{\text{rel}} \sigma_{a+b \rightarrow c+d}, \quad (4)$$

where  $N$  is the number of candidate scatters per second per unit volume, and  $v_{\text{rel}}$  is the relative velocity of the initial-state particles  $a$  and  $b$ . In case of a two-particle bound state,  $N = |\psi(0)|^2$  is the probability that both particles are at the origin.