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NUMERICAL SIMULATION OF MHD PROCESSES IN THE MUON COLLIDER TARGET

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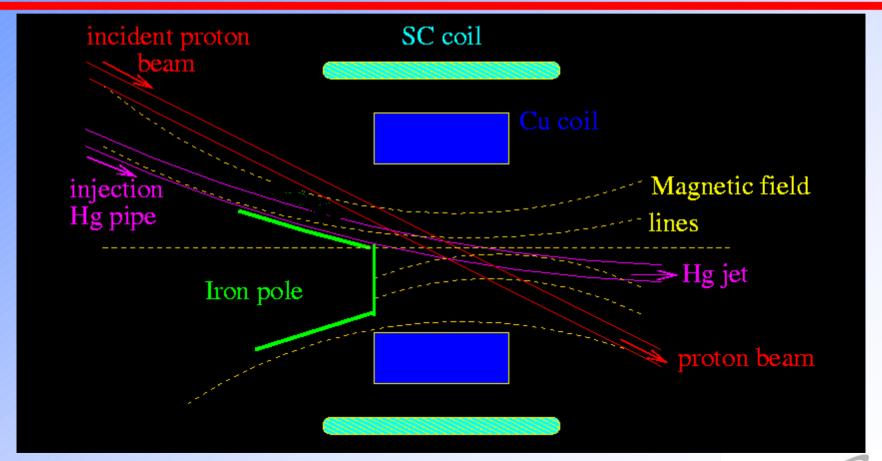


Talk outline

- FronTier-MHD code: theory and implementation.
- Numerical simulation of conducting liquid jets in magnetic fields and the muon collider target
- Modeling of the equation of state with phase transition (Riemann problem for such EOS)
- Numerical simulation of the mercury jet interacting with proton pulses.
- Plans for future research.



Schematic of the Muon Collider Target

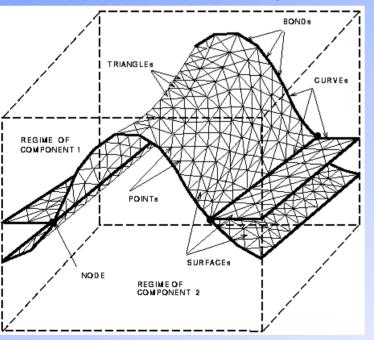


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3D FronTier Structures

Interfaces model different types of discontinuities in a medium such as shock waves in gas dynamics, boundaries between fluid-gas states, different fluids or their phases in fluid dynamics, component boundaries in solid dynamics etc.

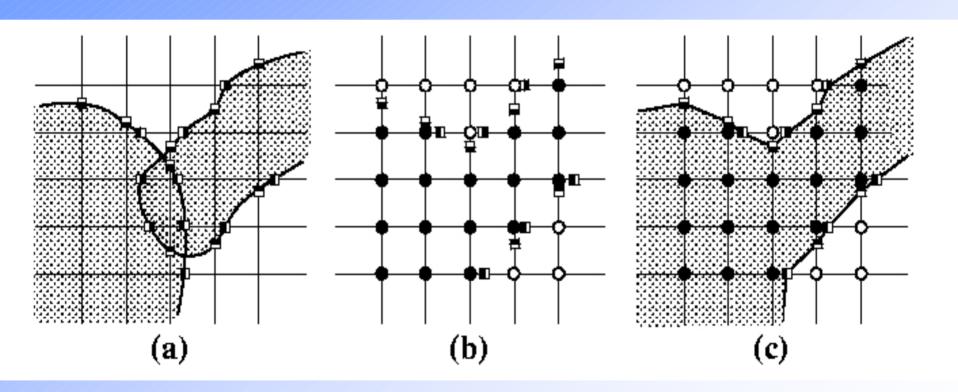


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Front tracking represents interfaces as lower dimensional meshes moving through a volume filling grid. The traditional volume filling finite difference grid supports smooth solutions located in the region between interfaces and the lower dimensional grid or interface defines the location of the discontinuity and the jump in the solution variables across it.



Resolving interface tangling by using the grid based method





The *FronTier* code: capabilities of the interior solvers

FronTier uses high resolution methods for the interior hyperbolic solvers such as Lax-Wendroff, Godunov and MUSCL and the following Riemann solvers:

- Exact Riemann solver
- Colella-Glaz approximate Riemann solver
- Linear US/UP fit (Dukowicz) Riemann solver
- Gamma law fit

FronTier uses realistic models for the equation of state:

- Polytropic Equation of State
- Stiffened Polytropic Equation of State
- Gruneisen Equation of State
- SESAME Tabular Equation of State

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The system of equation of compressible magnetohydrodynamics

$$\frac{\partial \rho}{\partial t} = -\nabla \cdot (\rho \mathbf{u})$$

$$\rho \left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) \mathbf{u} = -\nabla P + \rho \mathbf{X} + \frac{1}{c} (\mathbf{J} \times \mathbf{B})$$

$$\rho \left(\frac{\partial}{\partial t} + \mathbf{u} \cdot \nabla \right) U = -P \nabla \cdot \mathbf{u} + \frac{1}{\sigma} \mathbf{J}^2 - \frac{1}{c} \mathbf{u} \cdot (\mathbf{J} \times \mathbf{B})$$

$$\frac{\partial \mathbf{B}}{\partial t} = \nabla \times (\mathbf{u} \times \mathbf{B}) - \nabla \times \left(\frac{c^2}{4\pi\sigma} \nabla \times \mathbf{B} \right)$$

$$\nabla \cdot \mathbf{B} = 0$$



The interface condition

Interface conditions for the normal and tangential components of the magnetic field

$$\mathbf{n} \cdot (\mathbf{B}_2 - \mathbf{B}_1) = 0,$$

$$\mathbf{n} \times (\mathbf{H}_2 - \mathbf{H}_1) = \frac{4\pi}{c} \mathbf{K},$$
where $\mathbf{B} = \mu \mathbf{H}.$



Approximations: I. General case

General approach: the magnetic field is not constant in time.

Time scales: acoustic waves time scale = 7 microseconds;

magnetic diffusion time scale = 33 microseconds

Alfven waves time scale = 70 microseconds



Approximations: II. Constant in time magnetic field

The magnetic field is constant in time. The distribution of currents can be found by solving Poisson's equation:

$$\mathbf{J} = \sigma \left(\nabla \phi + \frac{1}{c} \mathbf{u} \times \mathbf{B} \right)$$

$$\Delta \phi = -\frac{1}{c} \nabla \cdot (\mathbf{u} \times \mathbf{B}),$$
with $\frac{\partial \phi}{\partial \mathbf{n}} \Big|_{\Gamma} = -\frac{1}{c} (\mathbf{u} \times \mathbf{B}) \cdot \mathbf{n}$



Magnetohydrodynamics of Multi Fluid Systems: Numerical Approach

- The system of MHD equations contains the hyperbolic subsystem (the mass, momentum and energy conservation equations) and the parabolic (the magnetic field evolution equation) or elliptic (Poisson's equation for the current density distribution) subsystems.
- The hyperbolic subsystem is solved on a finite difference grid in both domains separated by the free surface using FronTier's interface tracking numerical techniques. The evolution of the free fluid surface is obtained through the solution of the Riemann problem for compressible fluids.
- The parabolic subsystem or elliptic subsystems is solved using a vector finite elements method based on Whitney elements. The grid is rebuilt at every time step and conformed to the evolving interface.



Finite element mesh generation with interface constraints: Point shift method.

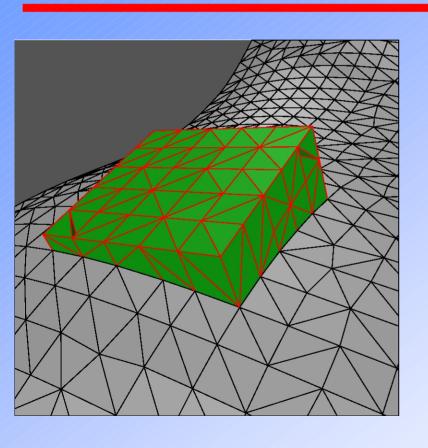
- Distortion of appropriate mesh blocks into irregular, convex hexahedra by shifting appropriate corner nodes to be coincident with the triangulated surface.
- Appropriate redistribution of surface triangles to ensure that surface triangles are coplanar either with shifted hexahedra faces or "interior diagonal" planes.
- Tetrahedralization of all (regular and irregular) grid blocks, creating and modifying a regularly indexed grid to provide a restricted optimized match to a triangulated tracked surface.

Advantages: preserves rectangular index structure.

Disadvantages: not robust (at present) in parallel.



Finite element mesh generation with interface constraints

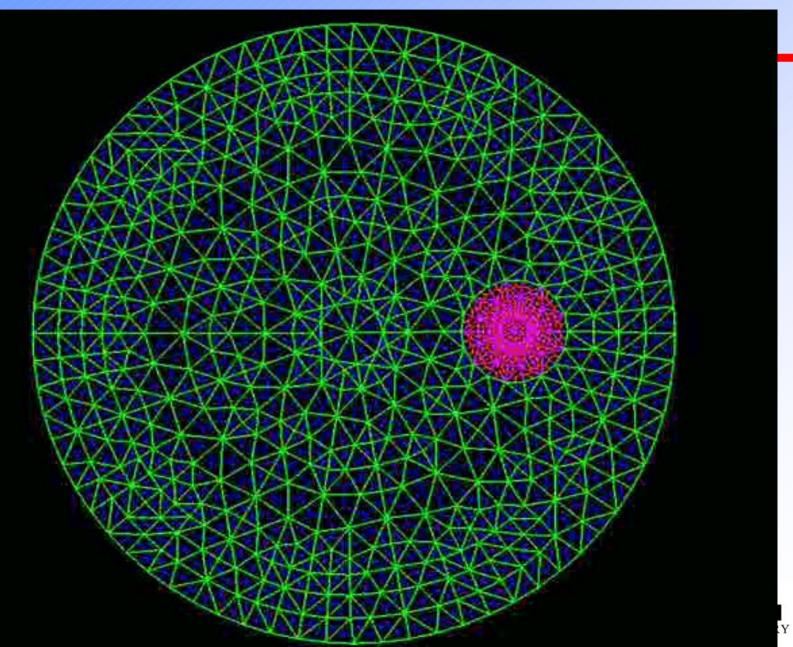


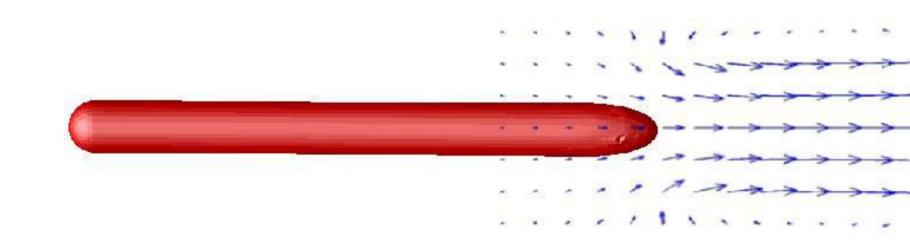
Triangulated tracked surface and tetrahedralized hexahedra conforming to the surface. For clarity, only a limited number of hexahedra have been displayed.

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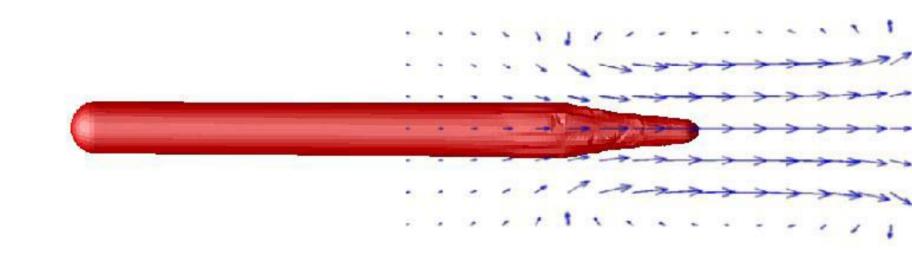


Dynamic finite element grid generation

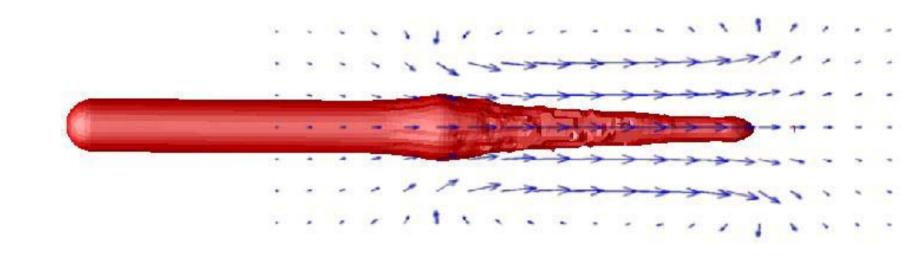




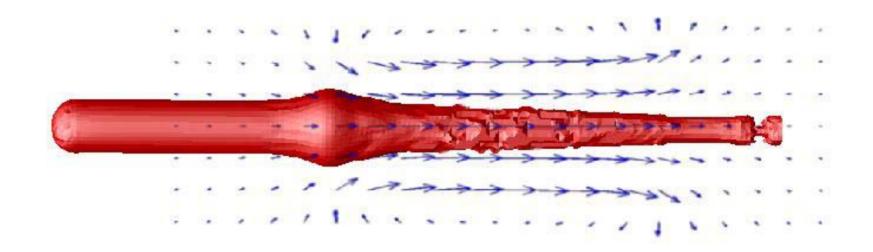




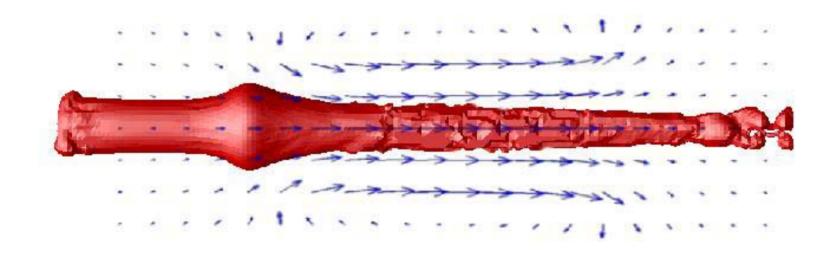






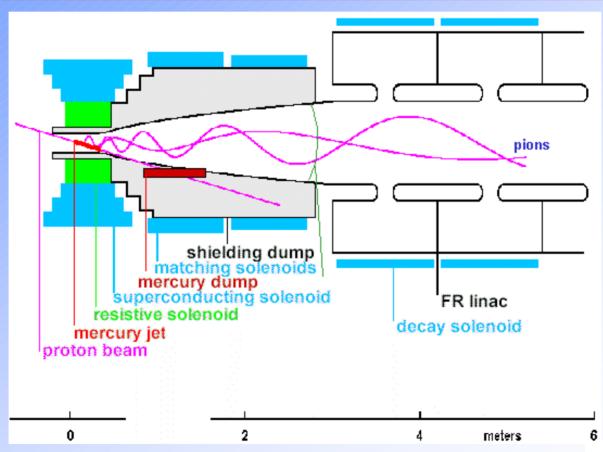






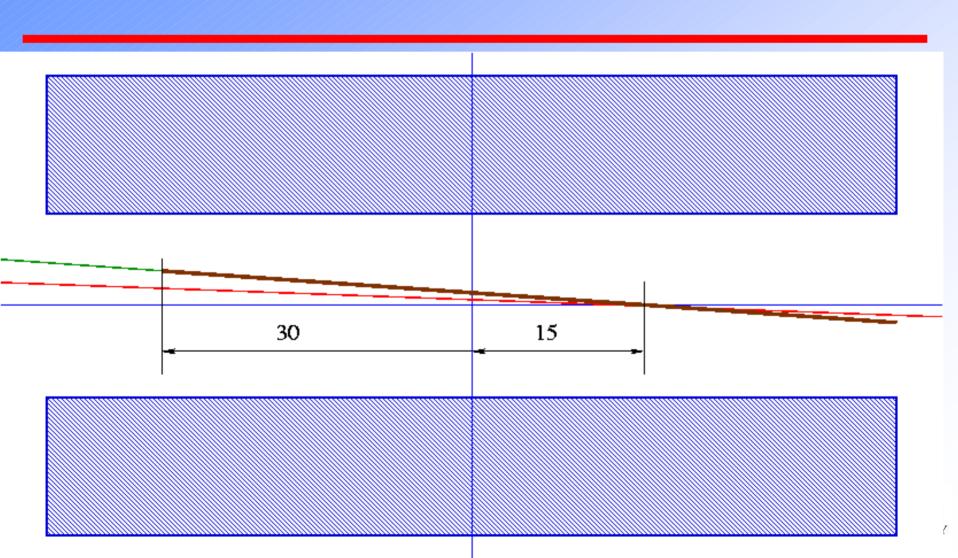


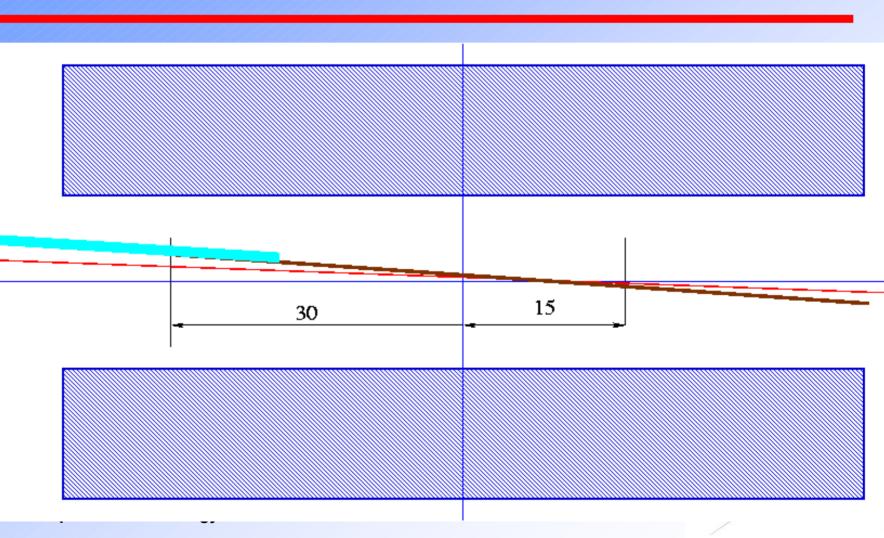
Schematic of the muon collider magnet system



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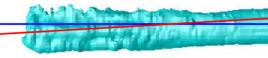








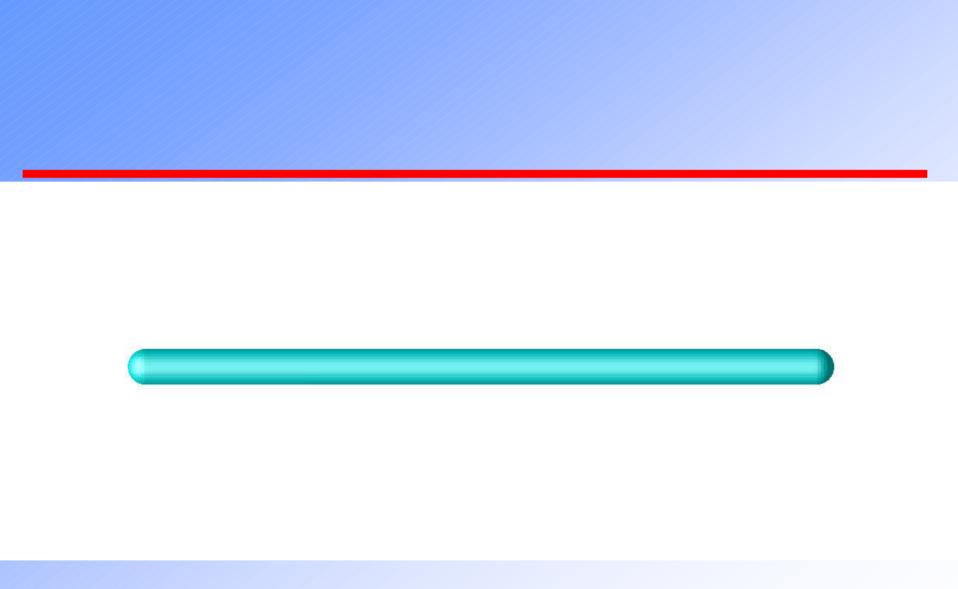




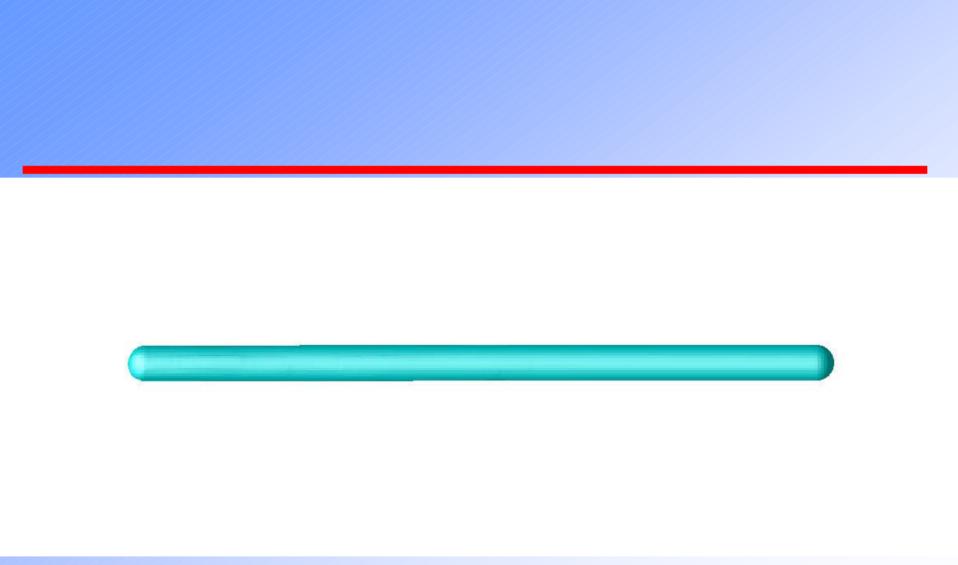


15 Tesla magnetic field

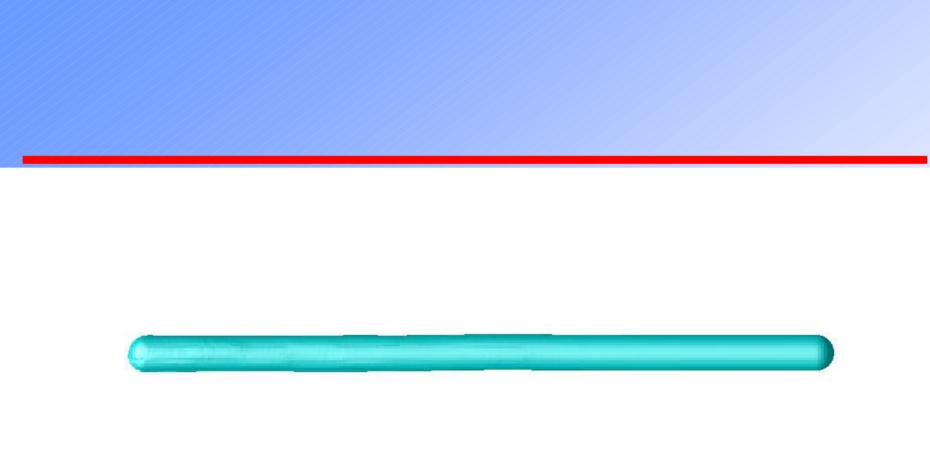




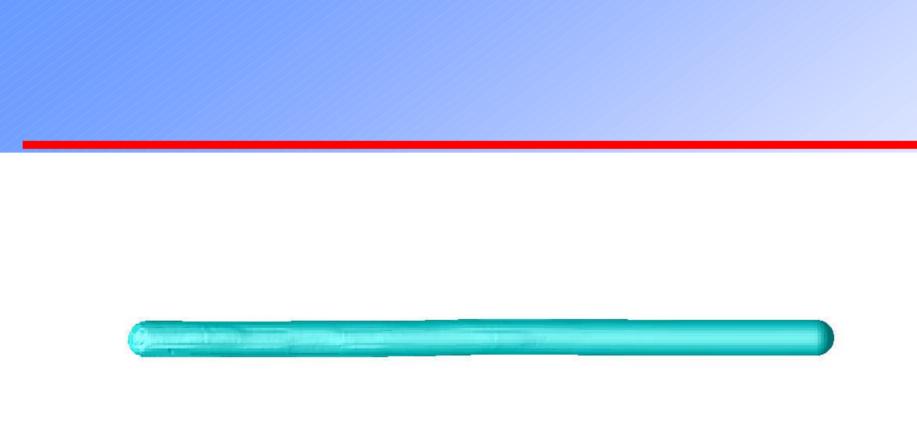




















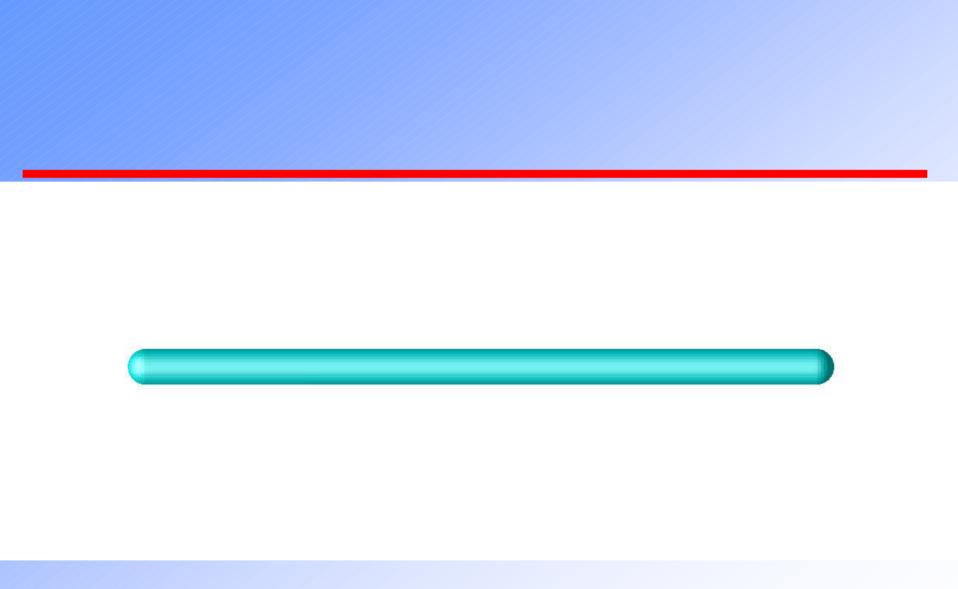




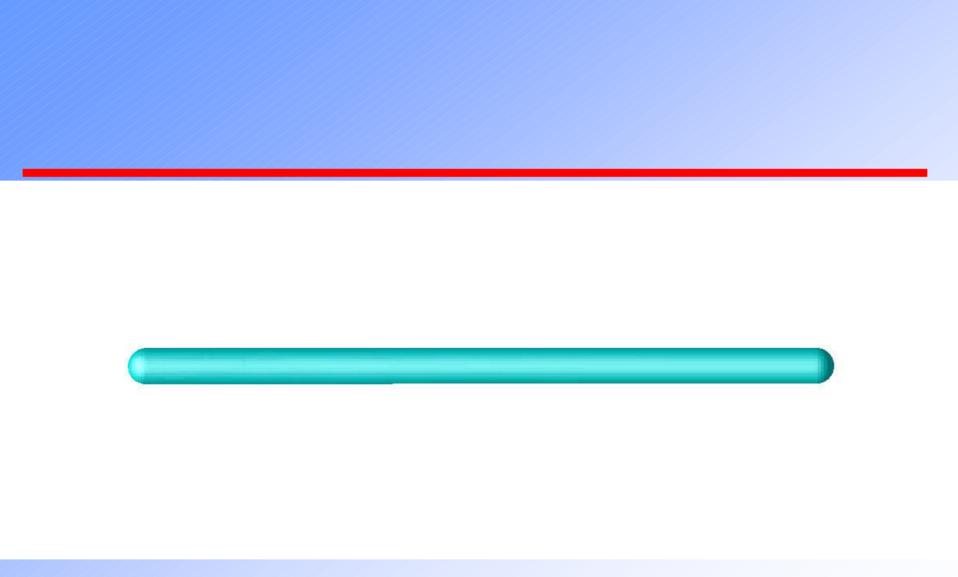


10 Tesla magnetic field

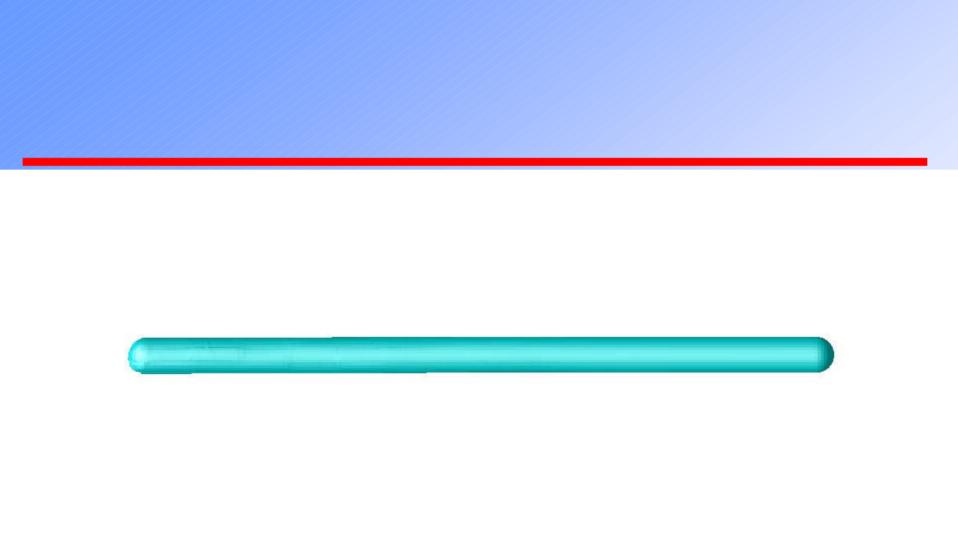




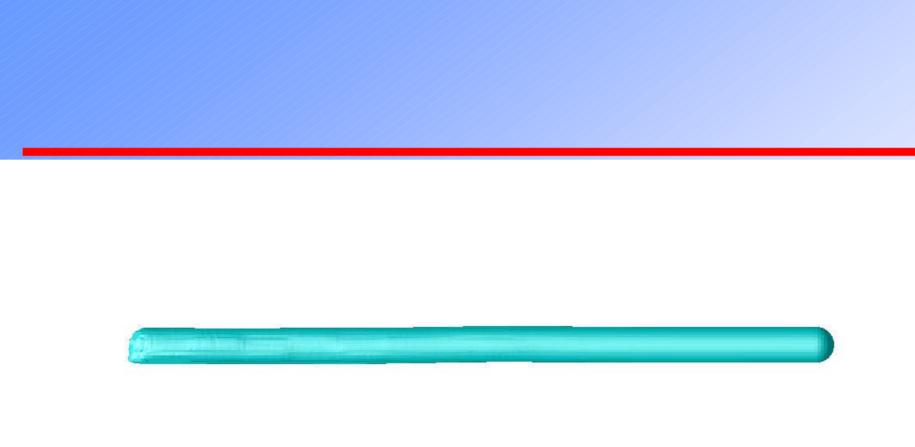




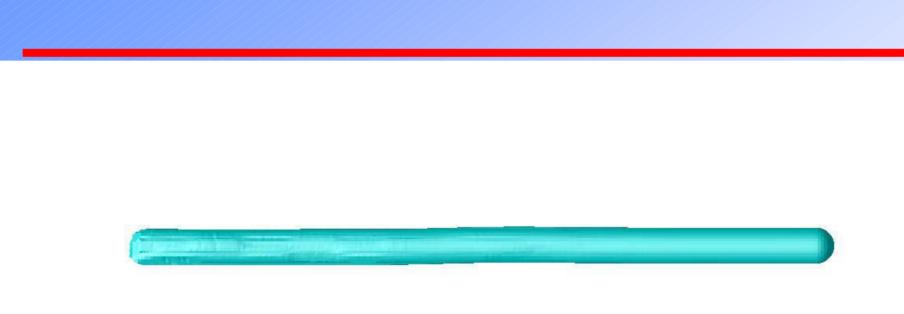














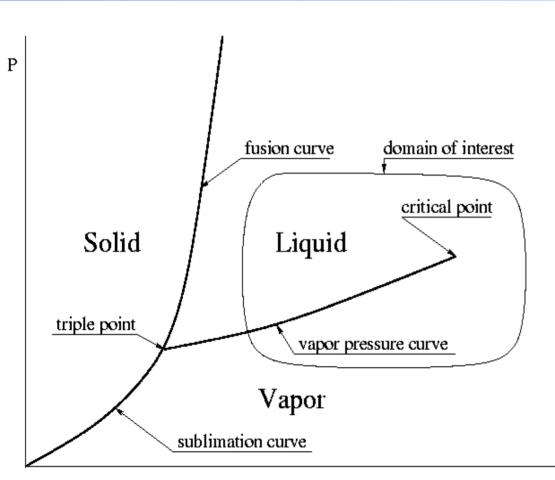








Mercury phase diagram



Critical point: $T_c = 1750$ K,

 $P_c = 172 \,\mathrm{MPa}, \ V_c = 43 \,\mathrm{cm}^3 \mathrm{mol}^{-1}$

Boiling point: $T_b = 629.84$ K,

$$P_b = 0.1 \text{MPa}, \rho = 13.546 \text{g} \cdot \text{cm}^{-3}$$

Analytical expression for the vapor pressure curve :

$$P = 133.3 \exp(18.41 - 7318 / T),$$

$$[P] = Pa, \quad [T] = K$$



Equation of state

An EOS is a relation expressing the specific internal energy E of a material as a function of the entropy S and the specific volume V: E(S, V).

The pressure P and temperature T are first derivatives of the energy E:

$$P(V,S) = -\frac{\partial E}{\partial V}\Big|_{S}, \qquad T(V,S) = \frac{\partial E}{\partial S}\Big|_{V}$$

in accordance with the second law of thermodynamics: TdS=dE+PdV.

The second derivatives of the internal energy are related to the adiabatic exponent γ , the Gruneisen coefficient Γ and the specific heat g.



SESAME tabular equation of state

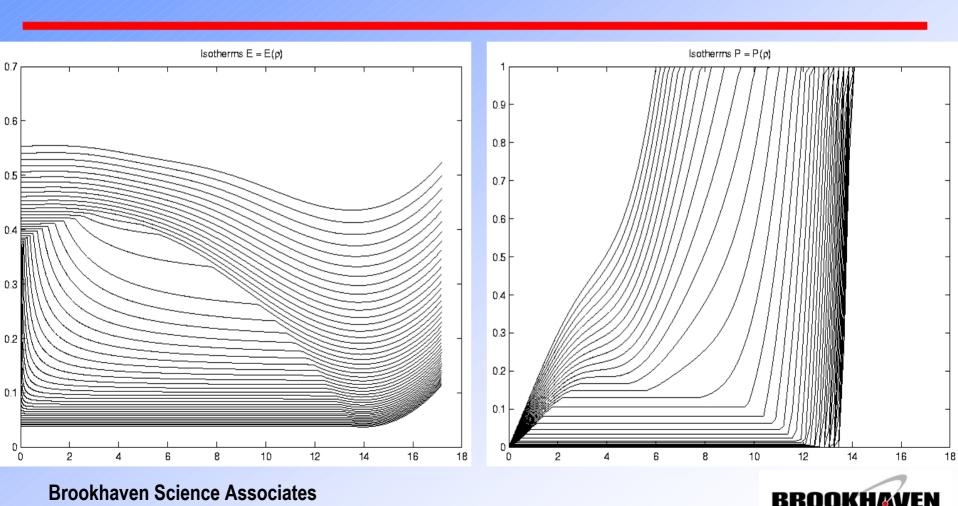
- The SESAME EOS is a tabular equation of state which gives the thermodynamic functions for a large number of materials, including gases, metals and minerals, in a computerized database.
- The SESAME EOS includes the following tables:
- **201 Tables.** The 201 tables contains 5 floats which are the *atomic number, atomic weight, density, pressure* and *internal energy* at the normal conditions.

301 Tables. The 301 tables is a database for the *pressure*, *internal energy* and, in some cases, the *free energy* as functions of the *temperature* and *density*.

401 Tables. The 401 tables contain data for the thermodynamic functions at the liquid/vapor phase transition.

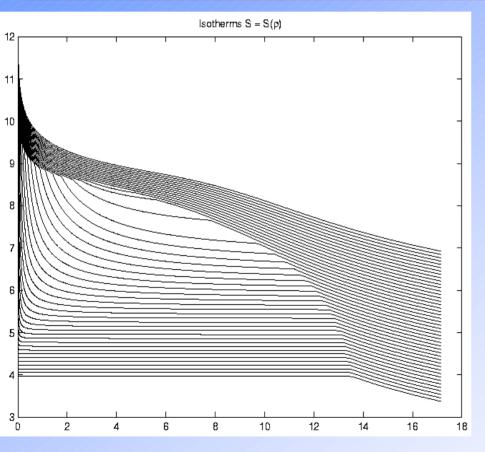
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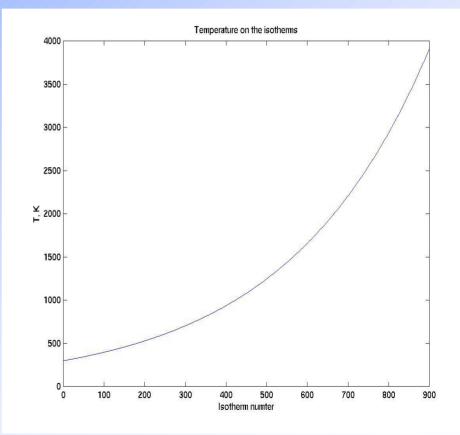
Thermodynamic properties of mercury, ANEOS data



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Thermodynamic properties of mercury, ANEOS data

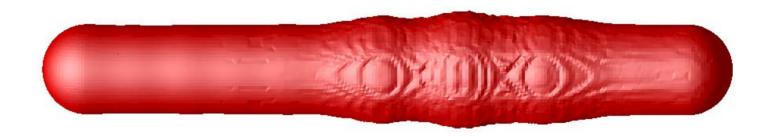




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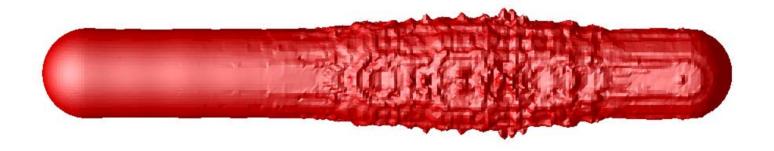


t = 27 mksec





t = 78 mksec





t = 120 mksec

