

# An R&D Program for Targetry and Capture at a Neutrino Factory and Muon Collider Source

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## Abstract

The need for intense muon beams for muon colliders and for neutrino factories based on muon storage rings leads to a concept of 1-4 MW proton beams incident on a moving target that is inside a 20-T solenoid magnet, with a mercury jet as a preferred example. Novel technical issues for such a system include disruption of the mercury jet by the proton beam and distortion of the jet on entering the solenoid, as well as more conventional issues of materials lifetime and handling of activated materials in an intense radiation environment. As part of the R&D program of the Neutrino Factory and Muon Collider Collaboration, an R&D effort related to

targetry is being performed within the context of experiment E951 at Brookhaven National Laboratory, first results of which are reported here.

*Key words:* target, neutrino factory, muon collider

*PACS:* 29.25-t

## 1 THE TARGETRY CONCEPT

A muon collider [1] or a neutrino factory based on a muon storage ring [2–4] require intense beams of muons, which are obtained from the decay of pions produced in proton-nucleus collisions. To maximize the yield, pions of momentum near 300 MeV/c should be captured [7–9], as illustrated in Fig. 1. For proton energies above 10 GeV, the pion yield per unit of proton beam energy is larger for a high- $Z$  target [8], as shown in Fig. 2. For proton beam energies in the MW range, beam heating would melt or crack a stationary high- $Z$  target [10], so a moving target must be used. A mercury jet target is the main focus of BNL E951 [6], although R&D is also being conducted on a carbon target option [2,11,12] as might be suitable for a low-energy proton source [13], and conceptual studies have been carried out for rotating-band targets [14,15], a tantalum/water target [16], and a liquid-lithium target [17].

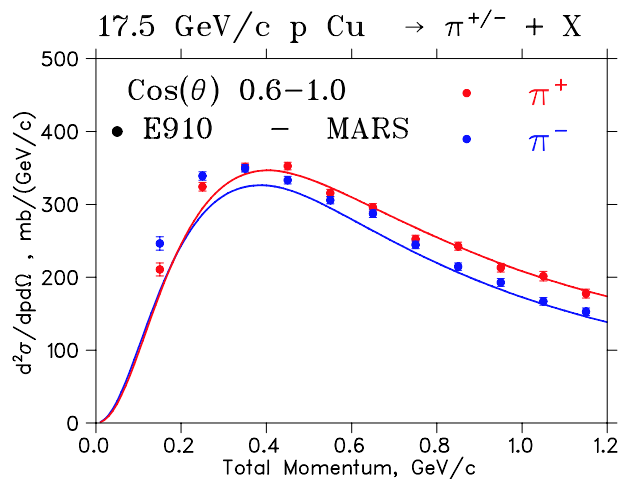


Fig. 1. Comparison of pion yield measured in BNL E910 [7] with a MARS calculation [8].

The low-energy pions are produced with relatively large angles to the proton beam, and efficient capture into a decay and phase rotation channel is obtained by surrounding the target with a 20-T solenoid magnet, whose field tapers

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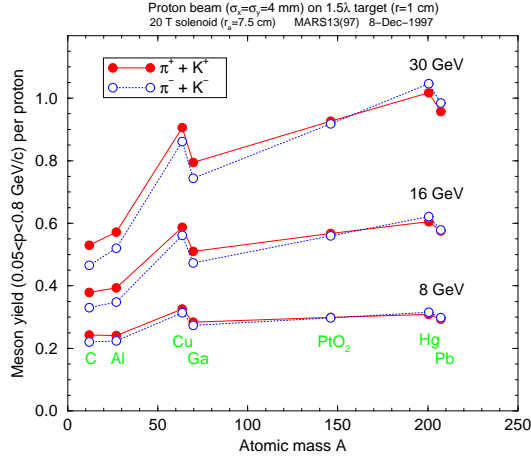


Fig. 2. Calculated yield of pion *vs.* atomic mass number of targets in proton beams of 8, 16 and 30 GeV [8].

down to 1.25 T over several meters [18,19], as sketched in Fig. 3. See also Figs. 2 and 3 of [9]. Pion yield is maximized with a mercury target in the form a 1-cm-diameter cylinder (Fig. 4), tilted by about 100 mrad with respect to the magnetic axis (Fig. 5). To permit the proton beam to interact with the target over 2 interaction lengths, the proton beam is tilted by 33 mrad with respect to the mercury jet axis.

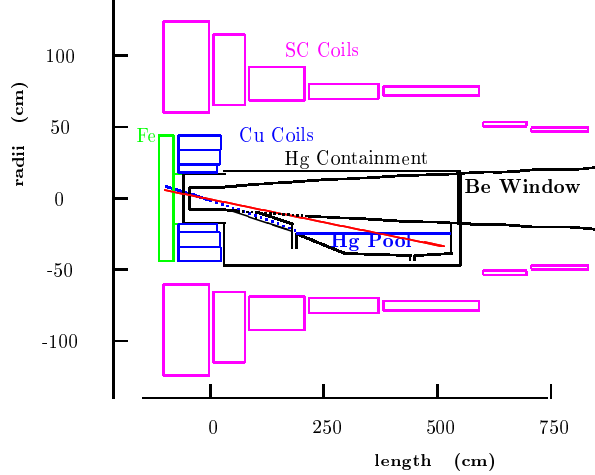


Fig. 3. Concept of targetry based on a mercury jet and proton beam at 100 mrad and 66 mrad, respectively, to the axis of a 20-T solenoid magnet.

The use of a mercury jet target raises several novel issues. The rapid energy deposition in the mercury target by the proton beam leads to intense pressure waves that can disperse the mercury [6,20–24]. Further, as the mercury enters the strong magnetic field eddy currents are induced in the mercury, and the Lorentz force on these currents could lead to distortion of the jet [6,24–29]. On the other hand, the magnetic pressure on the mercury once inside the solenoid will damp mechanical perturbation of the jet [21,30].

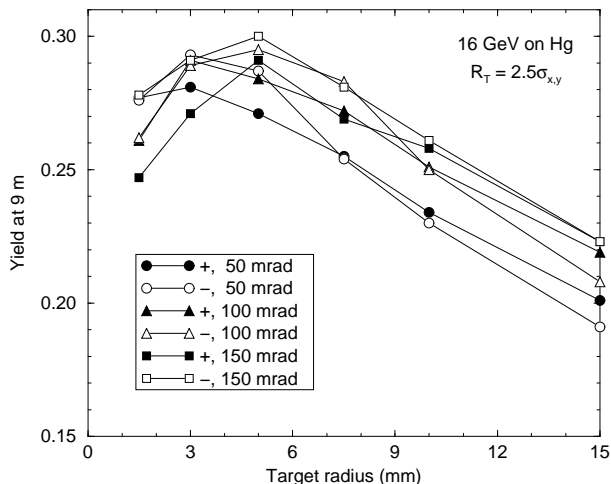


Fig. 4. Calculation of the yield of pions as a function of radius of a mercury target, assuming the proton beam has  $\sigma_r = 0.4r_{\text{target}}$  [8].

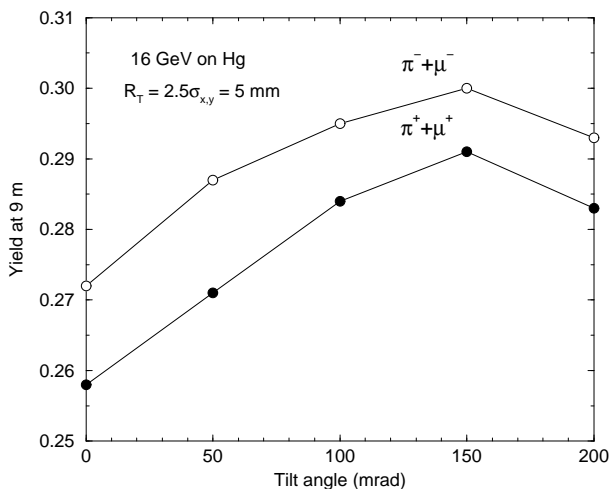


Fig. 5. Calculation of the yield of pions from a mercury target as a function of the tilt angle of the proton beam and target to the axis of a 20-T solenoid magnet [8].

To address these issues an R&D program is now underway.

## 2 THE TARGETRY R&D PROGRAM

In the USA, R&D on targetry for a neutrino factory and muon collider has been formalized as BNL experiment 951 [6]. This project maintains close contacts with related efforts in Europe [31] and in Japan [32].

The broad goal of E951 is to provide a facility that can test the major issues of a liquid or solid target in intense proton pulses and in a 20-T magnetic field. A sketch of the eventual facility is shown in Fig. 6.

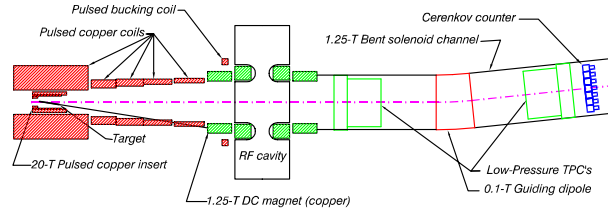


Fig. 6. Sketch of the full configuration of BNL E951, the targetry R&D facility.

Present activities in E951 focus on the interaction of intense proton pulses with targets in zero magnetic field. European targetry studies presently emphasize the interaction of mercury jets with a magnetic field, the operation of rf cavities near high-power targets [33], and evaluation of target materials [34].

### 2.1 Mercury Target Studies

The present R&D program on mercury jets is an outgrowth of work at CERN in the 1980's in which a prototype mercury jet was prepared (Fig. 7), but was never exposed to a beam.

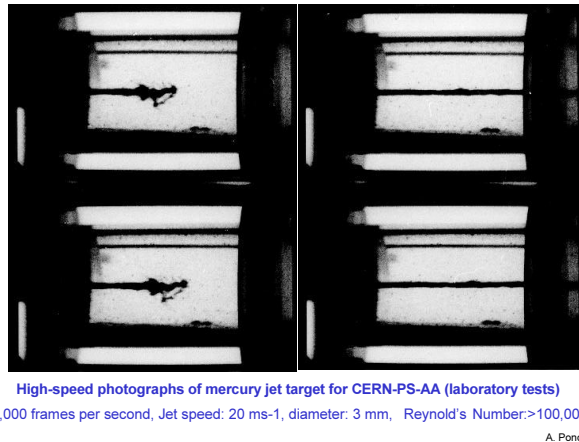


Fig. 7. Photographs of a 3-mm-diameter mercury jet (C.D. Johnson, 1988).

Experiment 951 is conducted in the A3 beamline of the BNL AGS [35] into which a single bunch, 100 ns long, of up to  $5 \times 10^{12}$  24-GeV protons can be extracted and brought to a focus as small as  $0.6 \times 1.6 \text{ mm}^2$ . The dispersal of both static and moving mercury targets by the proton beam was observed via two high-speed cameras using shadow photography with a laser diode [36]. Figure 8 shows the static target cell, and Figs. 9 and 10 show the effect of pulses of  $2\text{-}4 \times 10^{12}$  24-GeV proton on the mercury. Dispersal velocities of up to 50 m/s were observed. The air in the target cell slowed the droplet velocity by a factor of two over 10 cm. A key result from the jet studies was that the dispersal of mercury by the proton beam was confined to that part of the jet directly intercepted by protons.

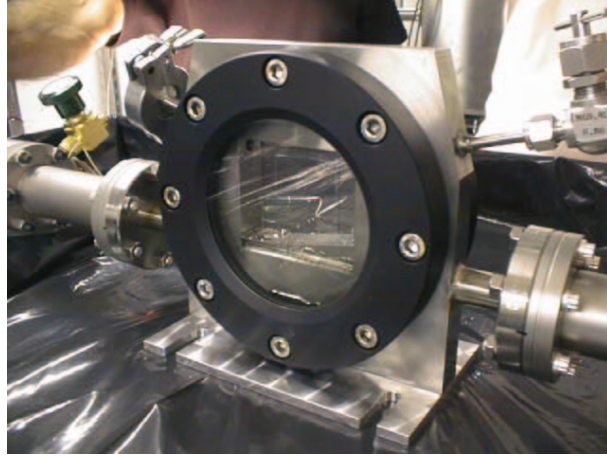


Fig. 8. Photograph of the static mercury target cell.

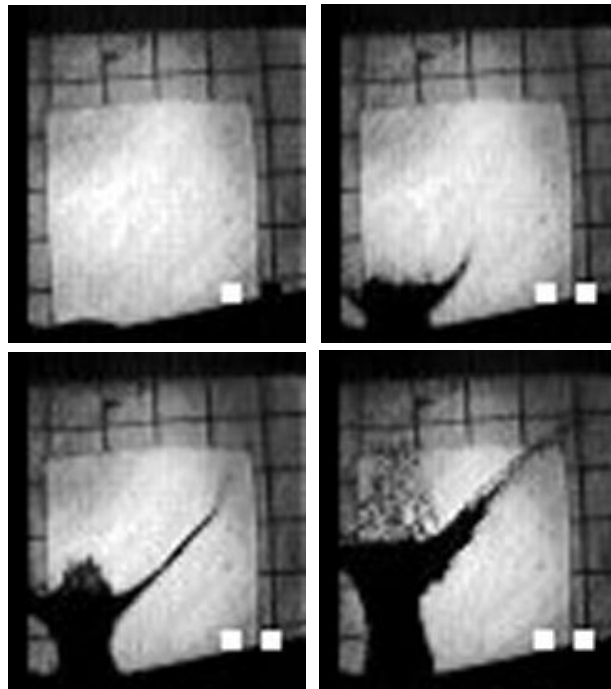


Fig. 9. Exposures of  $25 \mu\text{s}$  at  $t = 0, 0.5, 1.6, 3.4$  msec after a pulse of  $2 \times 10^{12}$  protons interacted with a static “thimble” of mercury 1.0 cm in diameter and 1.5 cm deep. The grid is  $1 \text{ cm} \times 1 \text{ cm}$ .

Thus, it appears that the dispersal of mercury by a proton beam is dramatic, but not violent, and that the dispersal will be a relatively modest issue for a target facility that operates at 15 Hz [39].

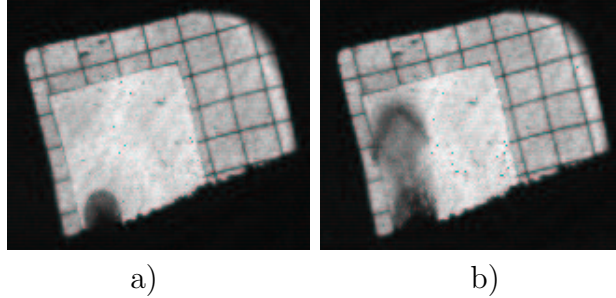


Fig. 10. Exposures of  $0.2 \mu\text{s}$  at  $t =$  a) 0.3, and b) 0.8 msec after a pulse of  $4 \times 10^{12}$  protons interacted with a static “thimble” of mercury. The grid is  $1 \text{ cm} \times 1 \text{ cm}$ .

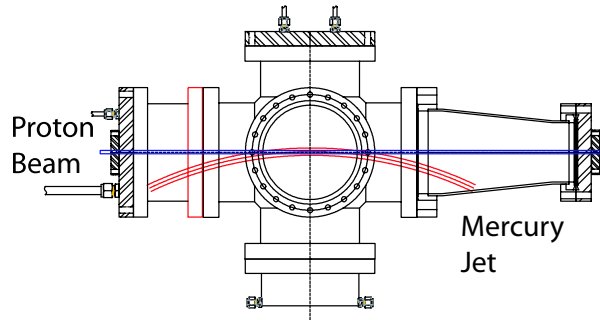


Fig. 11. Elevation view of the target cell that produced a 1-cm diameter jet of mercury with 2.5 m/s velocity that overlapped with the 24-GeV proton beam for 12 cm.

## 2.2 Solid Target Studies

E951 included exposures of several solid targets to the proton beam, using fiberoptic strain sensors with 500 kHz bandwidth (Fig. 13) to characterize the transient response of the targets to the pressure waves induced by beam energy deposition [40]. As expected, a carbon-carbon composite target with thermal expansion coefficient of less than  $10^{-6}/\text{K}$  showed much less strain than an ATJ graphite target (Fig. 15).

The issue of the rate of sublimation of carbon targets at the elevated temperatures ( $> 1900\text{C}$ ) caused by exposure to a 1-MW beam is under continuing laboratory study. Calculations indicate that a helium atmosphere can greatly extend the operational life of a carbon target against sublimation [41].

## 3 THE TARGET FACILITY

A preliminary design has been made of a target facility based on the concepts of the R&D program described above [3,9,42]. Low-energy pions are produced with relatively large angles to the proton beam, and efficient capture into a

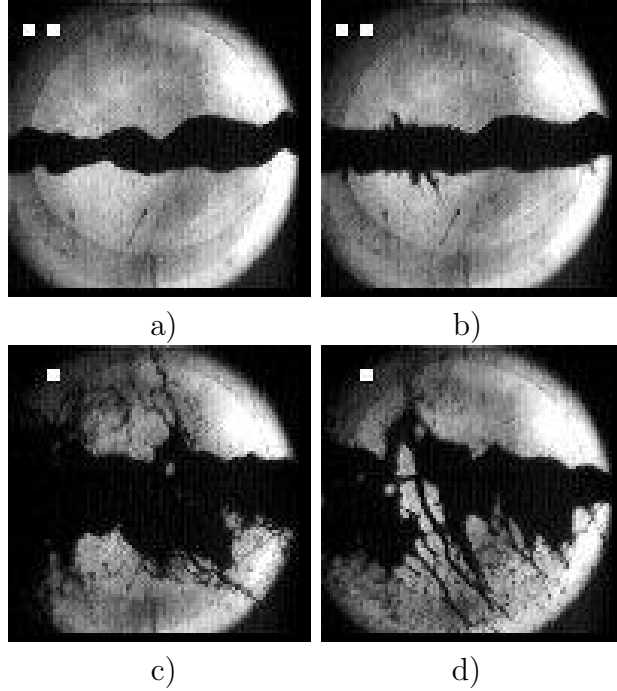


Fig. 12. Exposures of  $25 \mu\text{s}$  at  $t =$  a) 0, b) 0.75, c) 10, and d) 18 ms after a pulse of  $3.8 \times 10^{12}$  protons interacted with a free jet of mercury 1 cm in diameter.

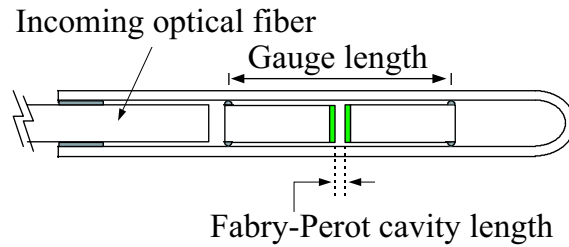


Fig. 13. Sketch of a fiberoptic strain sensor (FISO Technologies, <http://www.fiso.com>).

decay and phase rotation channel [43] is obtained by surrounding the target with a 20-T solenoid magnet, whose field tapers down to 1.25 T over several meters, as sketched in Fig. 16. Pion yield is maximized with a mercury target in the form a 1-cm-diameter cylinder, tilted by about 100 mrad with respect to the magnetic axis. To permit the proton beam to interact with the target over 2 interaction lengths, the proton beam is tilted by 33 mrad with respect to the mercury jet axis. See also Fig. 17.

A mercury pool inside the capture solenoid intercepts the mercury jet and the unscattered proton beam, as shown in Fig. 18. To suppress splashes of the mercury by the jet and/or proton beam as they enter the pool, a set of baffles can be arrayed above the pool and a particle bed submersed within it. The mercury pool, surrounding tungsten carbide/water shielding, and the



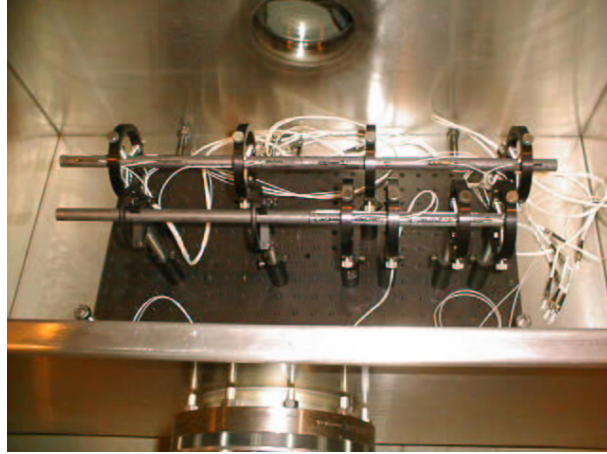


Fig. 14. View of the ATJ carbon and carbon-carbon composite targets equipped with fiberoptic strain sensors.

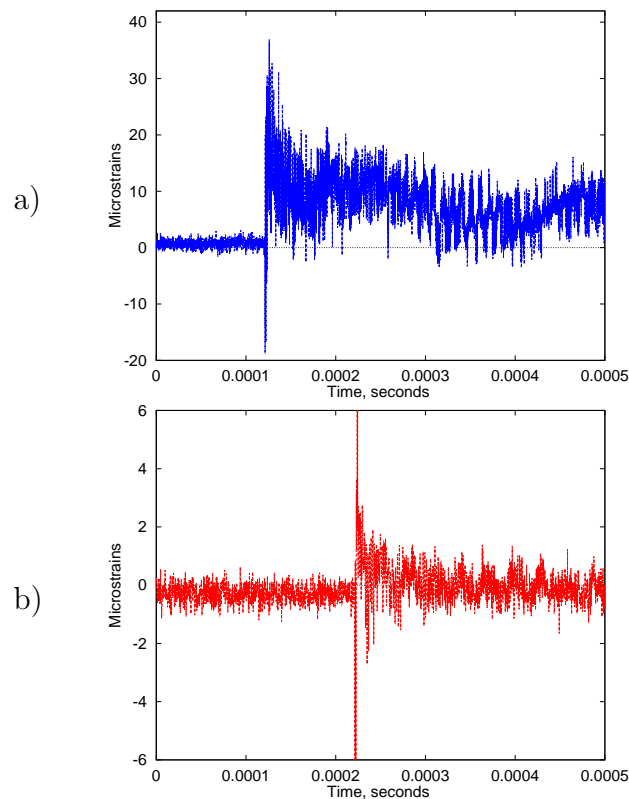


Fig. 15. Fiberoptic strain gauge data from carbon targets: a) a pulse of  $3 \times 10^{12}$  24-GeV protons on ATJ carbon; b) a pulse of  $2 \times 10^{12}$  protons on a carbon-carbon composite.

resistive insert of the 20-T capture magnet [19] are isolated from upstream and downstream beamline elements by a pair of double-walled Be windows. This entire unit can be replaced by remote manipulation should failure occur.

The absorbed radiation dose on components near the target is quite large [8],

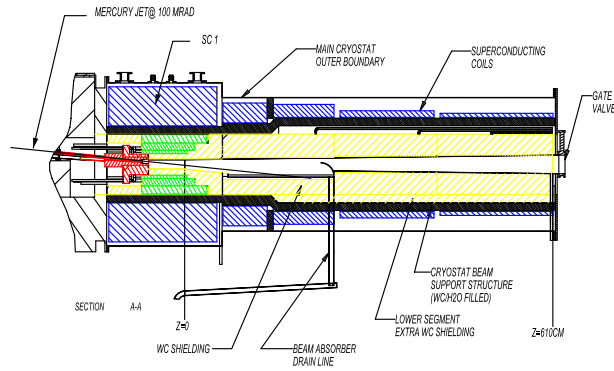


Fig. 16. Sketch of the target and capture system based on a mercury jet inside a 20-T solenoid magnet.

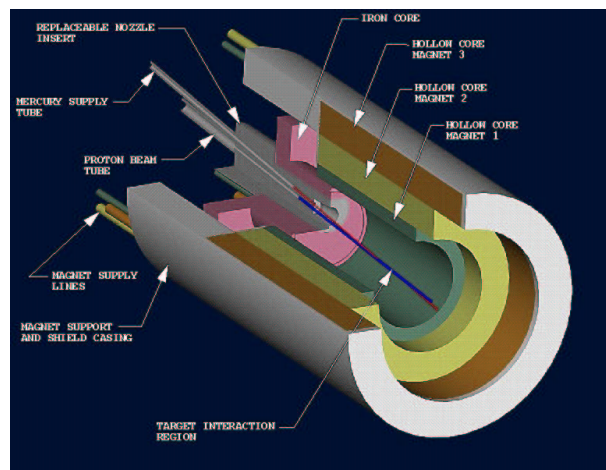


Fig. 17. The inner region of the 20-T capture magnet along with the tilted mercury jet target and proton beam.

as illustrated in Fig. 19, such that in a 4 MW proton beam their lifetime against radiation damage may only be 5 years.

The capture solenoid is encased in thick concrete shielding as part of the target facility that includes an overhead crane, hot cells with remote manipulation capability, and a mercury pumping and purification loop [42], as sketched in Fig. 20.

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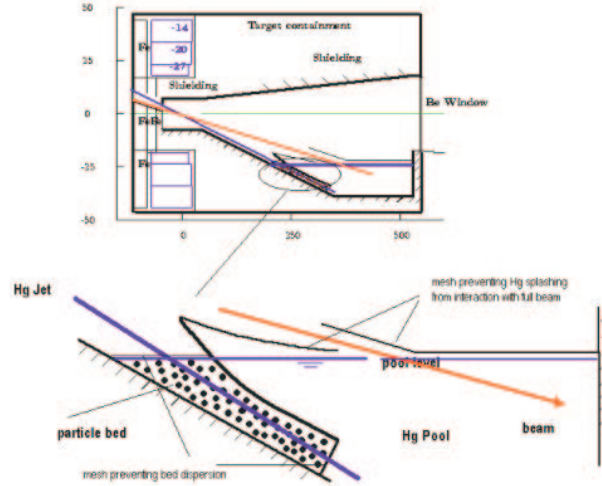


Fig. 18. A mercury pool inside the capture solenoid intercepts the mercury jet and the unscattered proton beam.

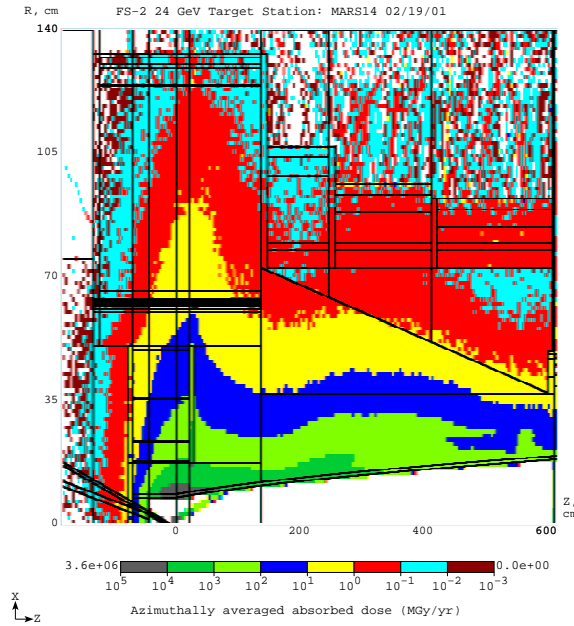


Fig. 19. Absorber radiation dose in the pion capture system per  $2 \times 10^7$  s of a 1-MW, 24-GeV proton beam on a mercury target.

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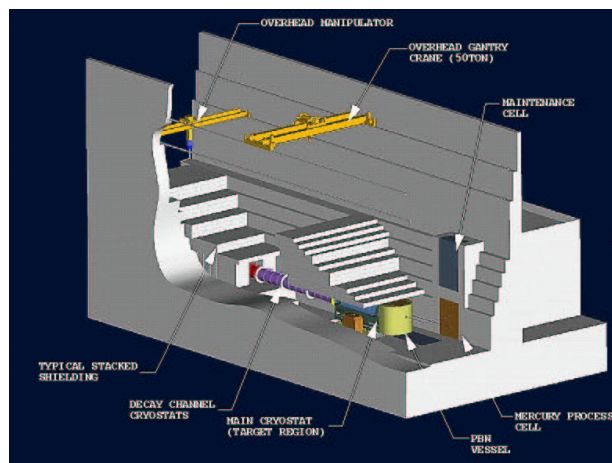


Fig. 20. Sketch of the target facility.

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