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Neutrino Factory / Muon Collider Targetry Meeting May 1 - 2, Oxford, GB

### **Target Simulations**

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### Talk Outline

![](_page_1_Picture_1.jpeg)

- Distortion of the mercury jet entering magnetic field
- Simulation of the mercury jet proton pulse interaction.
- Conclusions and future plans

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![](_page_1_Picture_8.jpeg)

### Main Ideas of Front Tracking

Front Tracking: A hybrid of Eulerian and Lagrangian methods

![](_page_2_Figure_2.jpeg)

#### Two separate grids to describe the solution:

- 1. A volume filling rectangular mesh
- 2. An unstructured codimension-1 Lagrangian mesh to represent interface

#### Major components:

- 1. Front propagation and redistribution
- 2. Wave (smooth region) solution

![](_page_2_Figure_9.jpeg)

- No numerical interfacial diffusion
- Real physics models for interface propagation
- Different physics / numerical approximations in domains separated by interfaces

![](_page_2_Figure_14.jpeg)

![](_page_2_Figure_15.jpeg)

### The *FronTier* Code

FronTier is a parallel 3D multiphysics code based on front tracking

- Physics models include
  - Compressible fluid dynamics
  - MHD
  - Flow in porous media
  - Elasto-plastic deformations
- Realistic EOS models
- Exact and approximate Riemann solvers
- Phase transition models
- Adaptive mesh refinement

![](_page_3_Figure_11.jpeg)

## Interface untangling by the grid based method

![](_page_3_Figure_13.jpeg)

### FronTier-MHD numerical scheme

![](_page_4_Figure_1.jpeg)

- or
- Perform finite volume discretization
- Solve linear system using fast Poisson solvers

### Main FronTier Applications

![](_page_5_Picture_1.jpeg)

Rayleigh-Taylor instability

#### Richtmyer-Meshkov instability

![](_page_5_Picture_4.jpeg)

![](_page_5_Picture_5.jpeg)

Tokamak refueling through the ablation of frozen D<sub>2</sub> pellets

![](_page_5_Picture_7.jpeg)

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### Jet entering 15 T solenoid

FronTier code:

- Explicitly tracked material interfaces
- Multiphase models
- MHD in low magnetic Reynolds number approximation

![](_page_6_Picture_6.jpeg)

![](_page_6_Picture_7.jpeg)

![](_page_7_Picture_0.jpeg)

### Previous Results (2005) Aspect ratio of the jet cross-section. I

![](_page_7_Figure_2.jpeg)

![](_page_7_Picture_5.jpeg)

![](_page_8_Picture_0.jpeg)

### Previous Results (2005) Aspect ratio of the jet cross-section. II

![](_page_8_Figure_2.jpeg)

![](_page_8_Picture_5.jpeg)

# Confirmation: Independent studies by Neil Morley, UCLA, HiMAG code

![](_page_9_Figure_1.jpeg)

![](_page_9_Figure_2.jpeg)

![](_page_9_Figure_3.jpeg)

![](_page_9_Figure_4.jpeg)

![](_page_9_Picture_5.jpeg)

![](_page_10_Picture_0.jpeg)

### Comparison with the theory

R. Samulyak et. al, Journal of Computational Physics, 226 (2007), 1532 - 1549.

![](_page_10_Figure_3.jpeg)

![](_page_10_Picture_4.jpeg)

![](_page_11_Picture_0.jpeg)

### **MERIT** setup

#### Geometry of Hg system in Magnet

![](_page_11_Figure_3.jpeg)

![](_page_11_Picture_4.jpeg)

![](_page_12_Picture_0.jpeg)

V = 15 m/s, B = 15 T

![](_page_12_Figure_2.jpeg)

![](_page_13_Picture_0.jpeg)

### V = 20 m/s, B = 15 T

![](_page_13_Figure_2.jpeg)

![](_page_14_Picture_0.jpeg)

### Comparison: V = 15 and 20 m/s, B = 10 and 15 T

![](_page_14_Figure_2.jpeg)

![](_page_14_Picture_3.jpeg)

![](_page_15_Picture_0.jpeg)

### **Experimental data**

V = 15 m/s, B = 10T

QuickTime<sup>™</sup> and a TIFF (LZW) decompressor are needed to see this picture.

B = 15T

QuickTime<sup>™</sup> and a TIFF (LZW) decompressor are needed to see this picture. V = 20 m/s, B = 10T

QuickTime<sup>™</sup> and a TIFF (LZW) decompressor are needed to see this picture.

B = 15T

QuickTime<sup>™</sup> and a TIFF (LZW) decompressor are needed to see this picture.

Simulations only qualitatively explain the width of the jet in different view ports.

![](_page_15_Picture_13.jpeg)

![](_page_16_Picture_0.jpeg)

Jet - proton pulse interaction. Evolution of models. Phase I: Single phase mercury (no cavitation)

- Strong surface instabilities and jet breakup observed in simulations
- Mercury is able to sustain very large tension
- Jet oscillates after the interaction and develops instabilities

#### Jet surface instabilities

![](_page_16_Figure_6.jpeg)

### Jet - proton pulse interaction. Phase II: Cavitation models

• We evaluated and compared homogeneous and heterogeneous cavitation models:

![](_page_17_Figure_2.jpeg)

![](_page_17_Figure_3.jpeg)

![](_page_17_Picture_4.jpeg)

Heterogeneous model (resolved cavitation bubbles)

- Two models agree reasonably well
- Predict correct jet expansion velocity
- Surface instabilities and jet breakup not present in in simulations

![](_page_17_Picture_9.jpeg)

![](_page_17_Picture_10.jpeg)

![](_page_17_Picture_11.jpeg)

### Jet - proton pulse interaction Phase II: Cavitation models in magnetic field

- The linear conductivity model predicts strong stabilizing effect of the magnetic field
- Stabilizing effect of the magnetic field is weaker if conductivity models with phase transitions are used (~ 20 % for Bruggeman's model)

![](_page_18_Figure_3.jpeg)

 If jet does not develop surface instabilities, the jet expansion is strongly damped in 15 T magnetic field (radial current are always present).
Experimentally confirmed in MERIT.

![](_page_18_Picture_7.jpeg)

![](_page_19_Picture_1.jpeg)

## Why surface instabilities and jet breakup are not observed in simulations with cavitation?

Possible Cause:

- Turbulence nature of the jet
- Microscopic mixture and strong sound speed reduction of the homogeneous model (separation of phases is important)
- Unresolved bubble collapse in the heterogeneous model
  - Bubble collapse is a singularity causing strong shock waves
- Other mechanisms?

![](_page_19_Picture_11.jpeg)

![](_page_20_Picture_0.jpeg)

### Multiscale approach to bubble collapse

• Bubble collapse (singularity) is difficult to resolve in global 3D model.

Multiscale approach:

Step 1: Accurate local model precomputes the collapse pressure

Step 2: Output of the local model serves as input to the global model

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![](_page_20_Picture_9.jpeg)

![](_page_21_Picture_0.jpeg)

### Step 1: 1D bubble collapse

Radius vs. Time

![](_page_21_Figure_3.jpeg)

#### Pressure Profile at t =0.0035 ms

Brookhaven Science Associates U.S. Department of Energy

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NA

![](_page_22_Figure_1.jpeg)

- Bubble collapse near the jet surface causes surface instability
- The growth of the spike is not stabilized by the magnetic field
- This is unlikely to be the only mechanism for surface instabilities

![](_page_22_Picture_7.jpeg)

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### Initial instability (turbulence) of the jet

![](_page_23_Picture_2.jpeg)

This was the real state of the jet before the interaction with protons

This was initial jet in previous numerical simulations (in 2D and 3D)

The obvious difference might be an important missing factor for both the jet flattening effect and interaction with the proton pulse.

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![](_page_23_Picture_8.jpeg)

![](_page_24_Picture_0.jpeg)

### 3D jet naturally growing from the nozzle

- Major numerical development allowed us to obtain the state of the target before the interaction by "first principles"
- Simulation of the jet proton pulse interaction is in progress

![](_page_24_Figure_4.jpeg)

![](_page_24_Figure_5.jpeg)

![](_page_24_Picture_6.jpeg)

### Mercury jet before the interaction with proton pulse. No magnetic field

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![](_page_25_Picture_2.jpeg)

![](_page_25_Picture_3.jpeg)