

## B(z) from Equations with Only One or Two Adjustable Currents

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Figure 1 plots on-axis field profiles generated by an equation, analytically differentiable to arbitrary order, each of whose terms predicts the on-axis field from a component coil in a Target Magnet. For simplicity, coils typically are modeled by current sheets. Comparison of the blue and black curves of Fig. reveals that even a coil as radially thick as Superconducting Coil #1 can be modeled accurately by a single current sheet at its mean radius.

For simplicity, each of the seven designs of Fig. 1 employs no more than three current elements, all those downstream of  $z \approx 5.4$  m being consolidated into a single current sheet. Design #1 uses only two elements—a main coil and a current sheet extending from 2.5 m all the way to 20 m. Coil parameters are tabulated below.  $B_\infty$  is the field that would be generated by a coil of identical current per unit length but of infinite length.

Design Number	#1	#2	#3	#4	#5	#6	#7
Desired field @ $z = L = 5$ m [T]	2.5	2.5	1.5	2.0	2.5	3.0	3.5
Inner radius $a_1$ of coil #1 [m]	1.20	1.20	N/A	N/A	N/A	N/A	N/A
Outer radius $a_2$ of coil #1 [m]	2.08	2.08	N/A	N/A	N/A	N/A	N/A
Mean radius $a_0$ of coil #1 [m]	1.64	1.64	1.64	1.64	1.64	1.64	1.64
$a_0$ of downstream coils [m]	1.00	1.00	1.00	1.00	1.00	1.00	1.00
Upstream end of coil #1 [m]	-2.35	-2.35	-2.35	-2.35	-2.35	-2.35	-2.35
Downstream end of coil #1 [m]	1.34	1.34	1.34	1.34	1.34	1.34	1.34
Upstream end of coil #2 [m]	2.50	2.50	2.50	2.50	2.50	2.50	2.50
Downstream end of coil #2 [m]	20.0	5.42	5.42	5.42	5.42	5.42	5.42
Upstream end of coil #3 [m]	N/A	5.427	5.427	5.427	5.427	5.427	5.427
Downstream end of coil #3 [m]	N/A	20.0	20.0	20.0	20.0	20.0	20.0
$j\lambda$ of thick-walled coil [ $A/mm^2$ ]	18.01	18.06	N/A	N/A	N/A	N/A	N/A
$B_\infty$ of upstream current sheet [T]	N/A	N/A	20.07	20.07	20.07	20.07	20.07
$B_\infty$ of middle current sheet [T]	N/A	1.684	0.623	1.153	1.684	2.214	2.744
$B_\infty$ of downstream sheet [T]	2.5	2.5	1.5	2.0	2.5	3.0	3.5

To satisfy constraints on the field magnitude  $B(z)$  and slope  $dB/dz$  at  $z = -0.5$  m and  $z = L$ , one can adjust the energization level of coils. Only coil currents—not coil dimensions—are variables, so that solving for variables requires no iteration. Furthermore, for the designs of Fig. 1 only the interior coil is used to satisfy the requirement that  $B(L) = B_{\min}$ . The upstream coil generates exactly 15 T without help for the downstream coils, which contribute only a few hundredth of a tesla at  $z = -0.5$  m. The downstream current sheet is energized so as to generate exactly  $B_{\min}$  deep in the interior of a coil of infinite length.

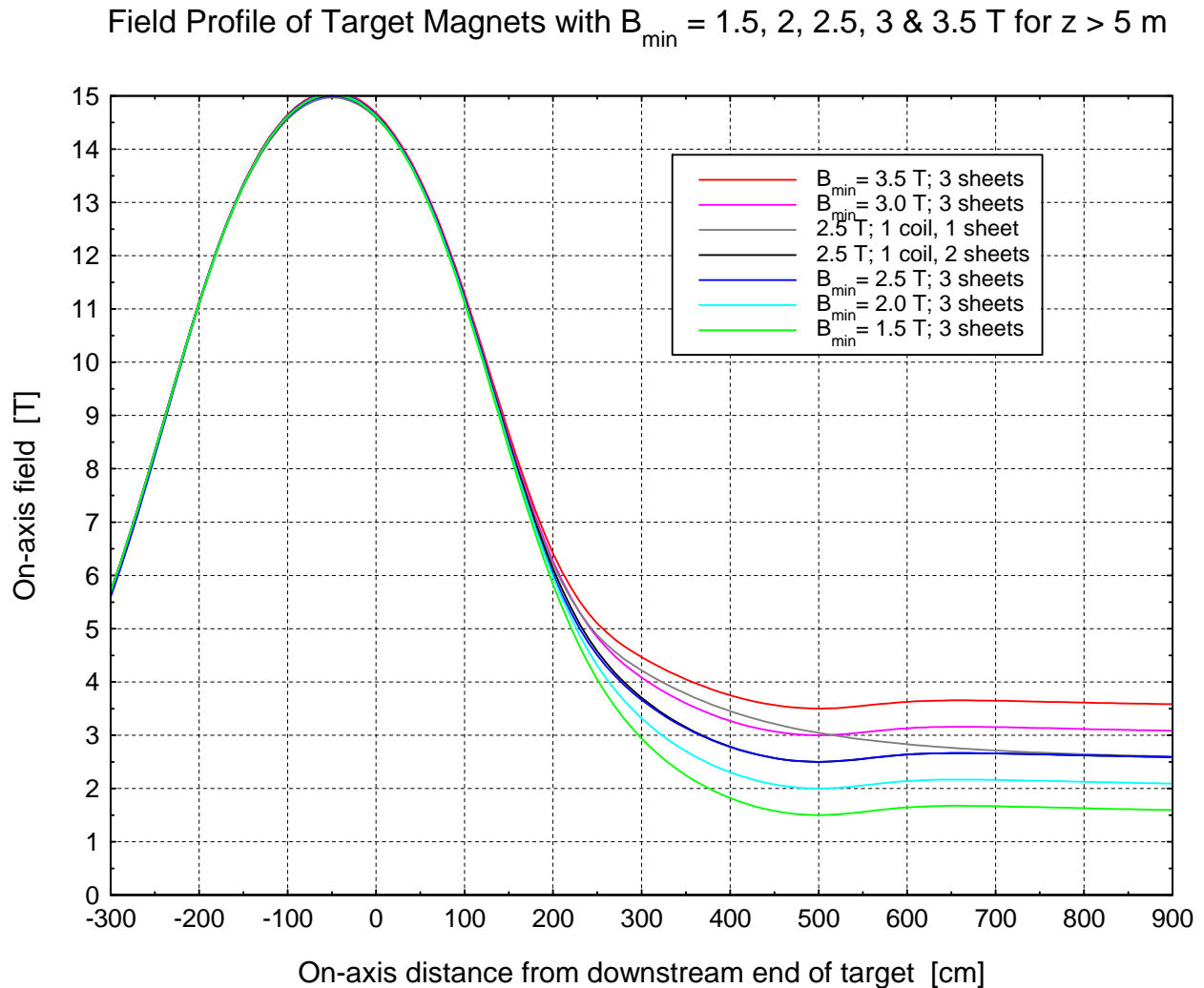


Fig. 1. On-axis field profiles generated by equations described in text above.

The field near  $z = 6$  m is higher than desired by as much as 0.17 T (0.5 T for the design with only one downstream coil). The field profiles of Figure 7 reduce this field deviation by means of an additional coil and therefore a second adjustable current. The solution of two simultaneous linear equations suffices to determine the currents in the interior two coils that satisfy the constraints  $B(z) = B_{\min}$  at  $z = 5$  m and at  $z = 6$  m. The latter constraint is more robust than requiring  $dB/dz = 0$  at  $z = 5$  m.

B(z) of Target Magnets with  $B_{\min} = 1.5, 2, 2.5, 3 \text{ \& } 3.5 \text{ T @ } z = 5 \text{ m}$

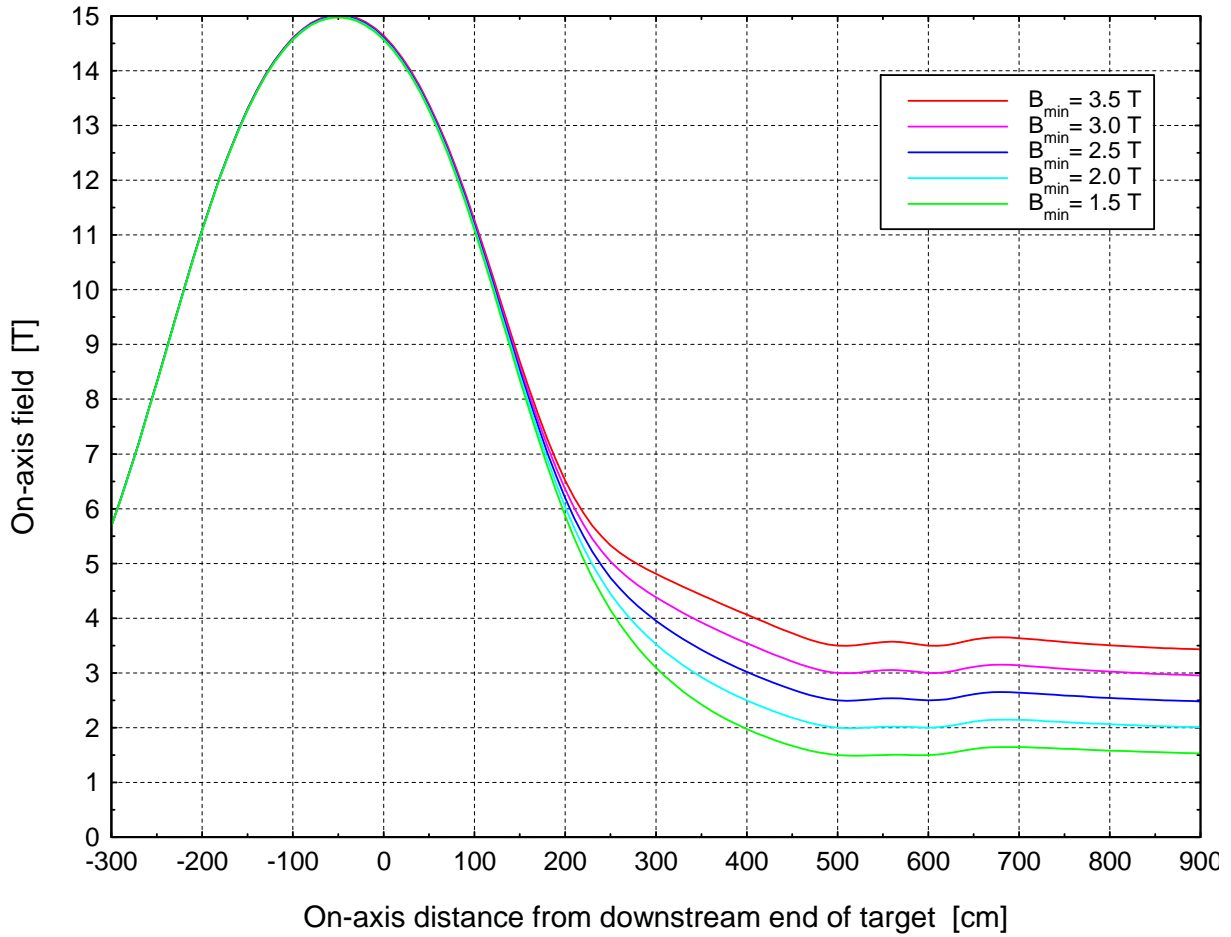


Fig. 2. On-axis field profiles generated by magnets with four current sheets. Currents  $i_1$  and  $i_4$  are fixed;  $i_2$  and  $i_3$  are adjustable.  $B(z) = B_{\min}$  at  $z = 5 \text{ m}$ ; also  $B(6\text{m}) = B_{\min}$ , a more robust constraint than  $dB/dz = 0$  at  $z = 5\text{m}$ .

All the designs employ four current sheets with consecutive radii of [1.64, 1.00, 0.75, 0.60] m; coil ends are at [-2.35, 2.50, 5.29, 6.28] m upstream and [1.34, 5.17, 6.18,  $\infty$ ] m downstream. To make the total field 15 T at  $z = -0.5\text{m}$ , the upstream coil generates 14.97 T. To make the total field approximately  $B_{\min}$  throughout the range 5 m to  $\sim 9 \text{ m}$ , the downstream coil carries 94% the current that it would need if unassisted by the coils upstream.

Design Number	#1	#2	#3	#4	#5
Desired field @ $z = L = 5 \text{ m}$ [T]	1.5	2.0	2.5	3.0	3.5
$B_{\infty}$ of current sheet #2 [T]	0.918	1.536	2.153	2.770	3.387
$B_{\infty}$ of current sheet #3 [T]	1.226	1.802	2.377	2.953	3.528

The on-axis field  $B$  at location  $z$  generated by a current sheet of radius  $a$ , upstream end  $b_1$ , downstream end  $b_2$ , and carrying  $i'$  amperes per meter is the function

$$B(z) = \frac{1}{2}\mu_0 i' (b-z) [a^2+(b-z^2)]^{-1/2} \text{ teslas}$$

evaluated at  $b = b_2$  minus its evaluation at  $b = b_1$ .

Define a function  $B^{(0)}(x) \equiv x(a^2+x^2)^{-1/2}$ , where  $x \equiv (b-z)$ . Similarly, define functions  $B^{(n)}(x) \equiv (1/n!) d^n B^{(0)}/dx^n$ . Then the Taylor series expansion of the field an increment  $\Delta$  from  $z$  is the function

$$\frac{1}{2}\mu_0 i' [B^{(0)}(x) + B^{(1)}(x) \Delta + B^{(2)}(x) \Delta^2 + B^{(3)}(x) \Delta^3 + B^{(4)}(x) \Delta^4 + B^{(5)}(x) \Delta^5 + \dots],$$

evaluated at  $x_1 \equiv (b_1-z)$  and subtracted from its evaluation at  $x_2 \equiv (b_2-z)$ .

Define  $r = (a^2+x^2)^{1/2}$ . The first five terms in the series are:

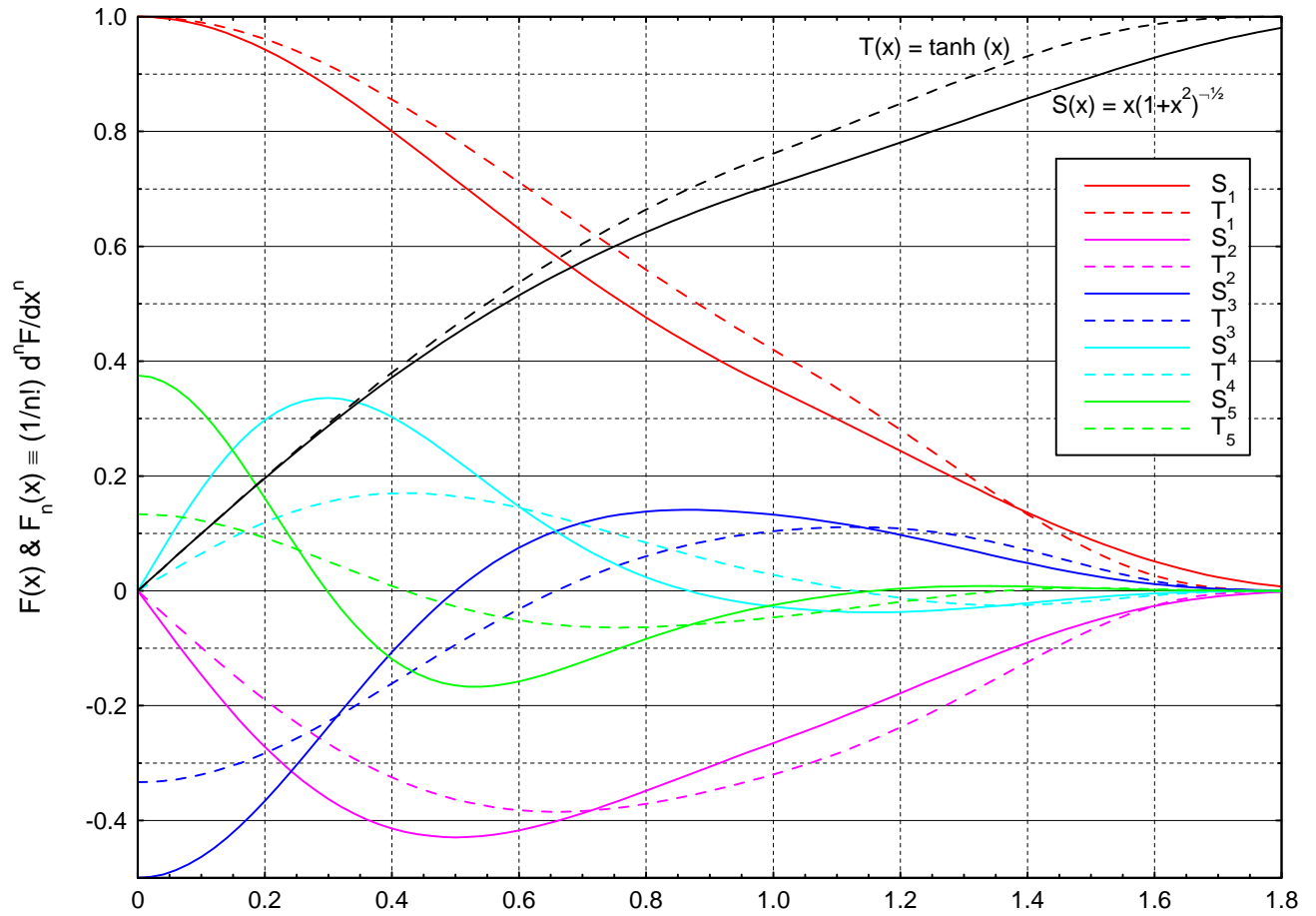
$$\begin{aligned} B^{(1)}(x) &= a^2 / r^3 , \\ B^{(2)}(x) &= -(3/2) a^2 x / r^5 , \\ B^{(3)}(x) &= -(1/2) (a^2 - 4 x^2) / r^7 , \\ B^{(4)}(x) &= (5/8) a^2 x (3 a^2 - 4 x^2) / r^9 , \\ B^{(5)}(x) &= (3/8) a^2 (a^4 - 12 a^2 x^2 + 8 x^4) / r^{11} . \end{aligned}$$

Figure 3 plots these functions with  $a = 1$ . To devote most of the graph to values of  $x$  less than unity, values greater than unity are transformed as  $x = (2-u)^{-1}$ , where  $u$  is the abscissa. The right-hand border of the graph,  $u = 1.8$ , transforms to  $x = 5$ .

The figure also plots the power-series coefficients for the function  $B^{(0)}(x) = \tanh(x)$ , which behaves much like  $(1+x^2)^{-1/2}$ . Its coefficients are:

$$\begin{aligned} B^{(1)}(x) &= 1 / \cosh^2(x) , \\ B^{(2)}(x) &= -\sinh(x) / \cosh^3(x) , \\ B^{(3)}(x) &= -(1/3) [1 - 2 \sinh^2(x)] / \cosh^4(x) , \\ B^{(4)}(x) &= (1/3) [\sinh(x)] [2 - \sinh^2(x)] / \cosh^5(x) , \\ B^{(5)}(x) &= (2/15) [2 - 11 \sinh^2(x) + 2 \sinh^4(x)] / \cosh^6(x) . \end{aligned}$$

### Derivatives of Functions $x/(1+x^2)^{1/2}$ and $\tanh(x)$



If  $x \leq 1$ ,  $u \equiv x$ . If  $x \geq 1$ ,  $u \equiv 2-x^{-1}$ ; Inverted,  $x \equiv (2-u)^{-1}$ . For example, if  $u = 1.8$ ,  $x = 5$ .

Fig. 3. Field and power-series coefficients of current sheet and function  $\tanh(x)$ , which behaves much like  $(1+x^2)^{-1/2}$ .